

# Quantum sensing with superconducting qubits

Logan Bishop-Van Horn  
QSQM Symposium  
2021-09-10

# Transmon Hamiltonian

PHYSICAL REVIEW A 76, 042319 (2007)

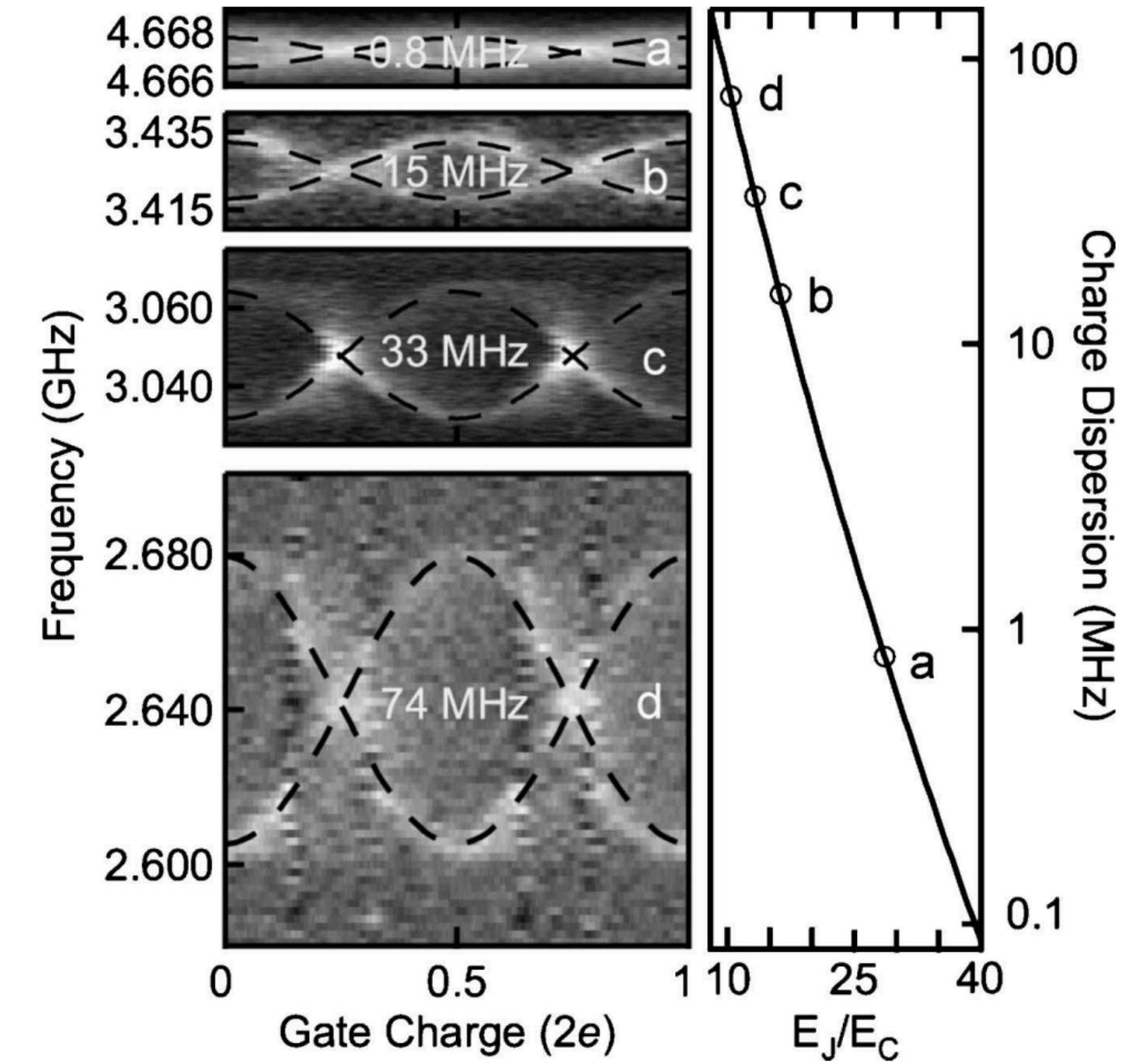
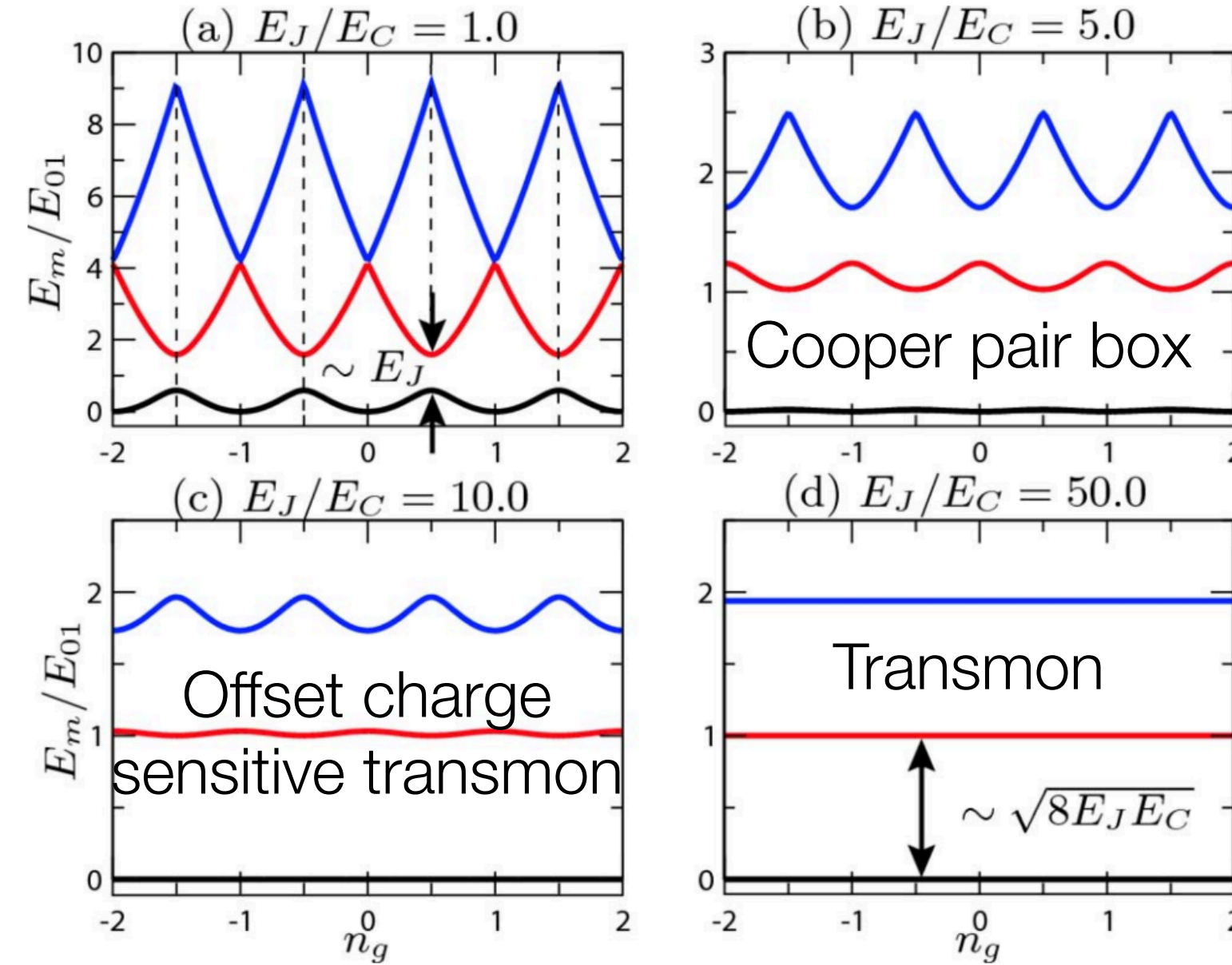
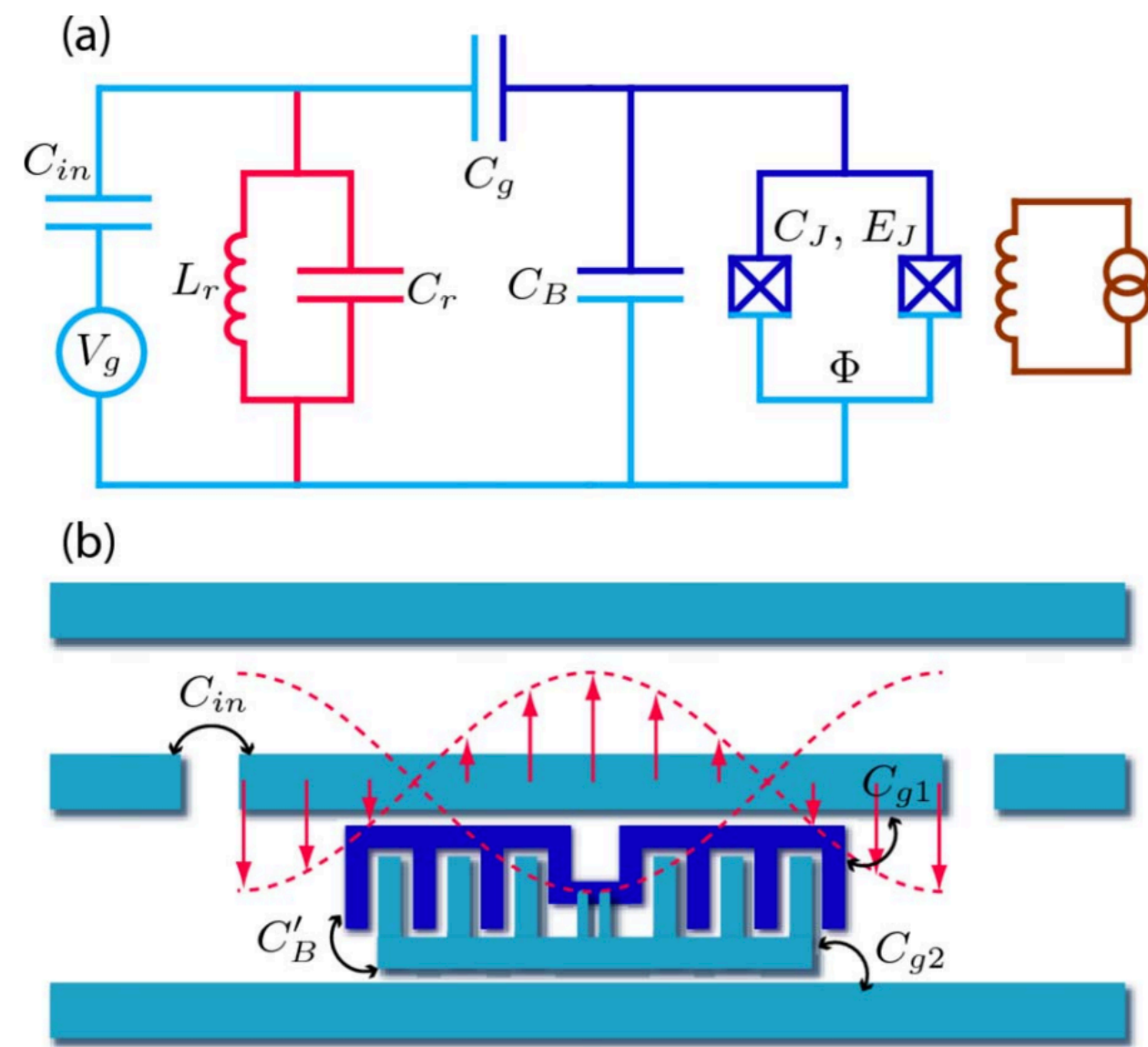
PHYSICAL REVIEW B 77, 180502(R) (2008)



Suppressing charge noise decoherence in superconducting charge qubits

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Charge-insensitive qubit design derived from the Cooper pair box



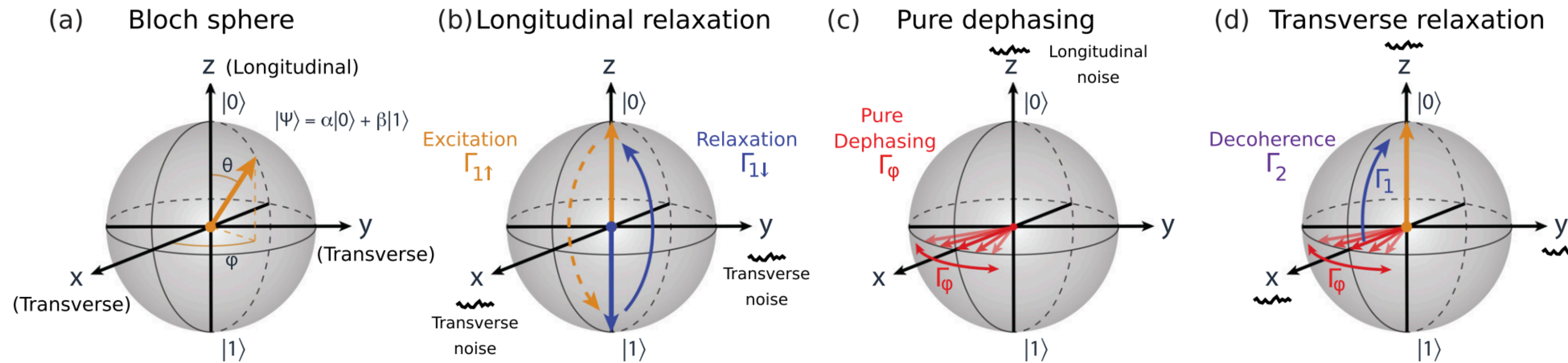
$$\hat{H} = 4E_C(\hat{n} - n_g)^2 - E_J(\Phi)\cos \hat{\phi}$$

- Superconducting circuit with tunable charge and flux sensitivity
- Transition frequencies  $\omega_{ij} = 2\pi f_{ij}$  are periodic in  $n_g$  and  $\Phi$
- $E_J/E_C$  determines anharmonicity and charge sensitivity



# Decoherence limits sensitivity

Cite as: Appl. Phys. Rev. 6, 021318 (2019); <https://doi.org/10.1063/1.5089550>



Goal: Measure qubit frequency  $f_{01}(n_g, \Phi)$  to quantify electric potential (via  $n_g$ ), magnetic field (via  $\Phi$ ), or fluctuations in those parameters (noise spectroscopy).

High frequency noise: drives transitions between states ( $T_1$ )

Control electronics  
Lossy environment  
Thermal radiation  
Pair-breaking photons

$$|\langle 0 | \hat{O}_\lambda | 1 \rangle|^2, S_\lambda(f_{01})$$

Low frequency (e.g.  $1/f$ ) noise: stochastically shifts qubit frequency ( $T_\varphi$ )

Control electronics  
Two-level systems  
Readout photons  
???

$$\frac{\partial f_{01}}{\partial \lambda}, S_\lambda(f)$$

$$T_2 = \left( \frac{1}{T_\varphi} + \frac{1}{2T_1} \right)^{-1}$$

$$\delta n_g \sim \frac{1}{T_2 \left| \frac{\partial f_{01}}{\partial n_g} \right|} \text{ Uncertainty in offset charge}$$

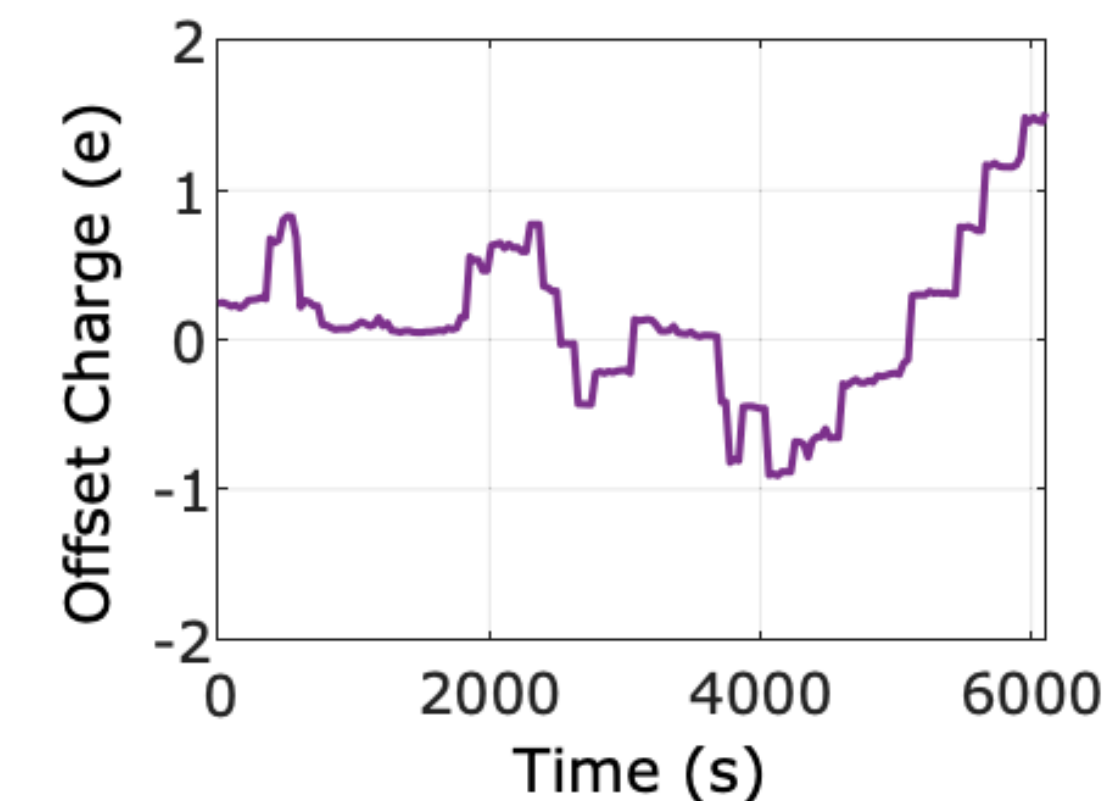
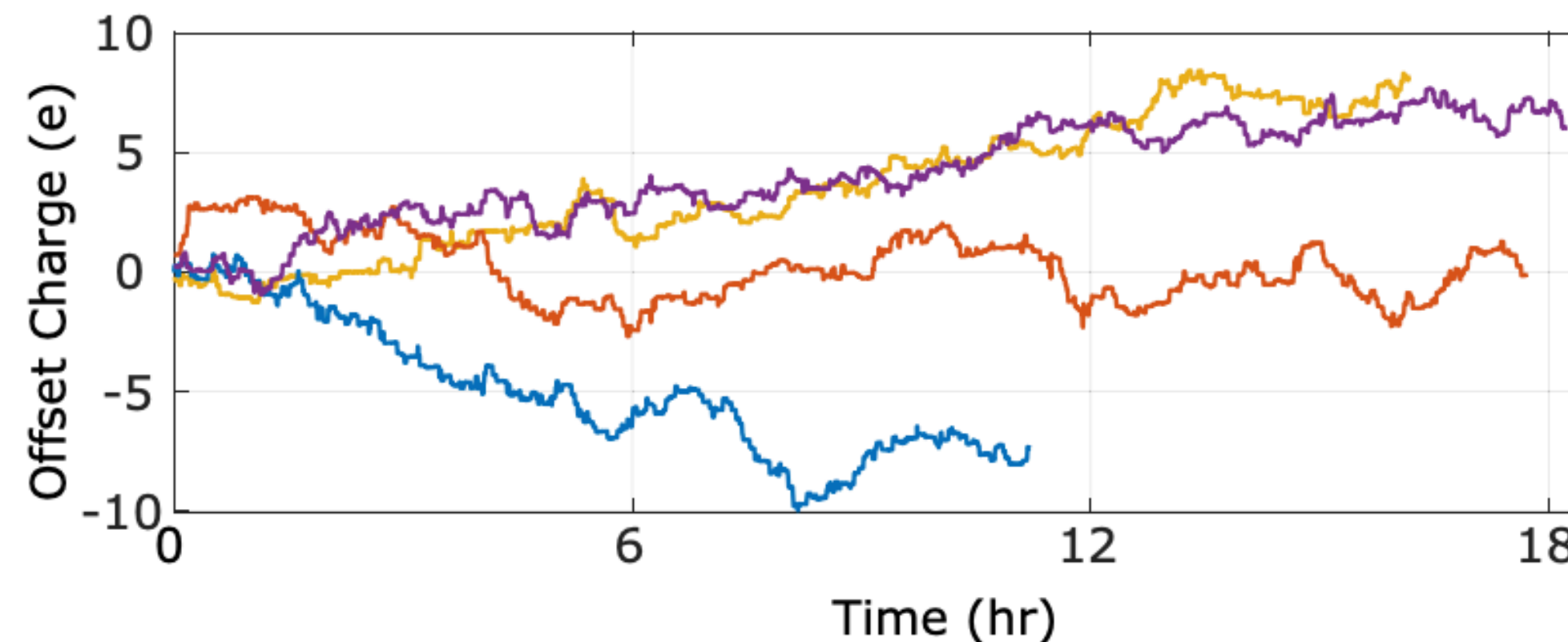
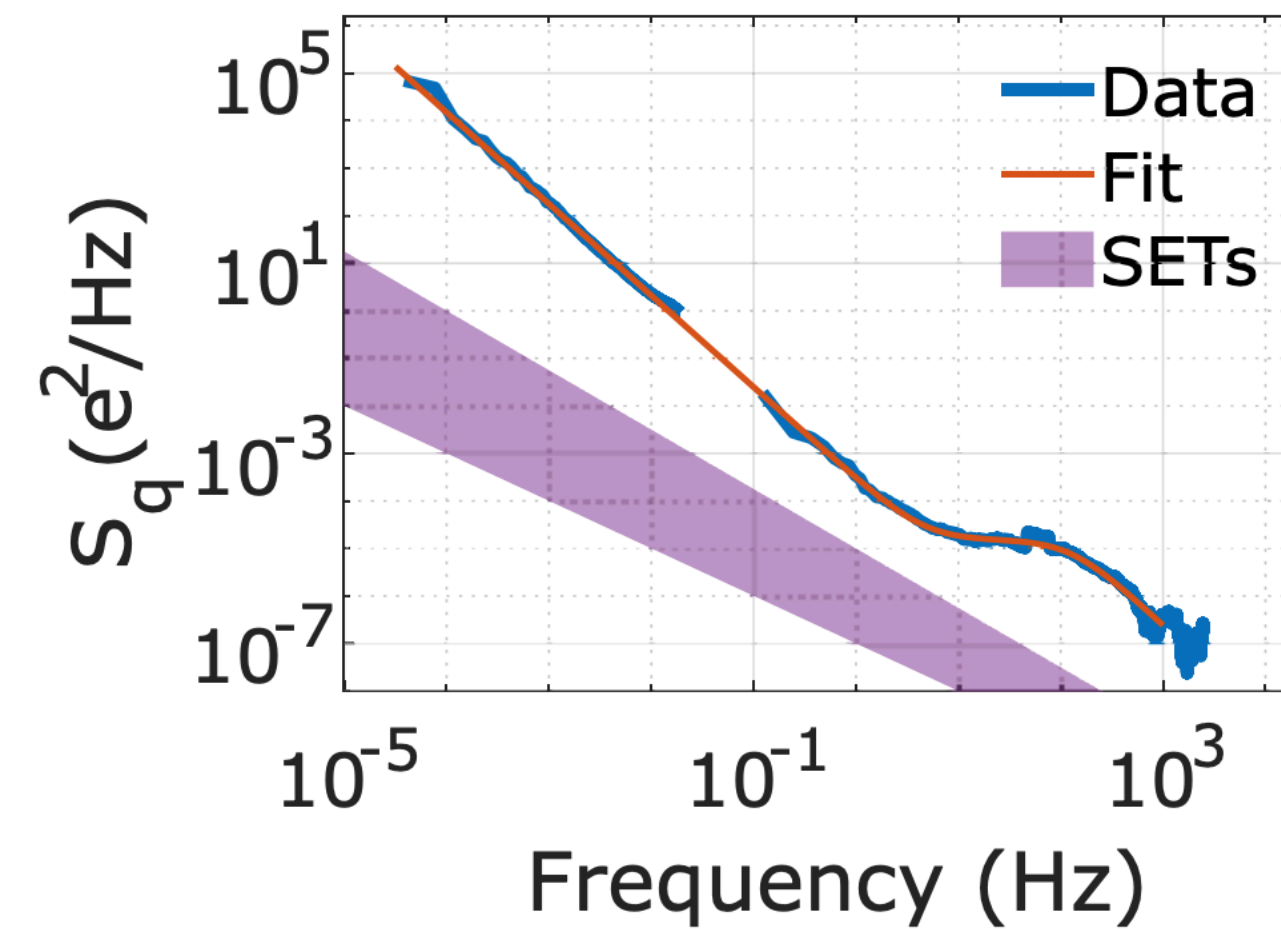
$$\delta \Phi \sim \frac{1}{T_2 \left| \frac{\partial f_{01}}{\partial \Phi} \right|} \text{ Uncertainty in flux bias}$$

# Low frequency charge noise

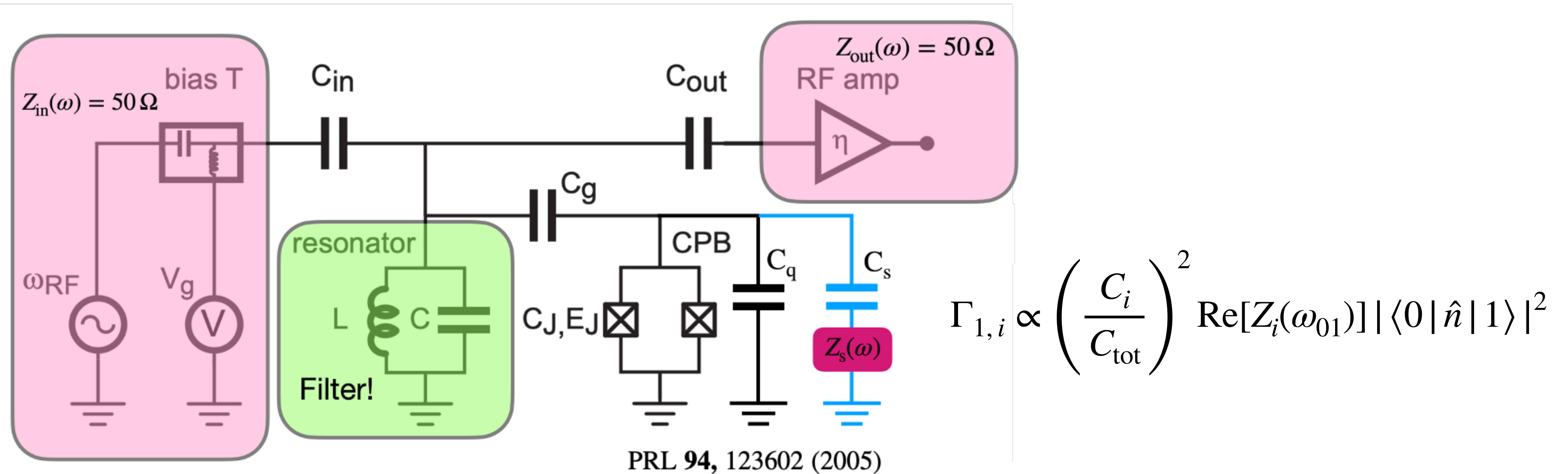
- 1/f-like charge noise in offset charge sensitive qubits is orders of magnitude worse than typically seen in SETs
- $\sim 0.1 - 1,000$  Hz: limits sensitivity via  $T_2$
- $\lesssim 0.1$  Hz: frequency drift complicates data-taking and analysis

## Anomalous charge noise in superconducting qubits

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# Coupling to sample with finite impedance



$$\Gamma_{1,i} \propto \left( \frac{C_i}{C_{\text{tot}}} \right)^2 \text{Re}[Z_i(\omega_{01})] |\langle 0 | \hat{n} | 1 \rangle|^2$$

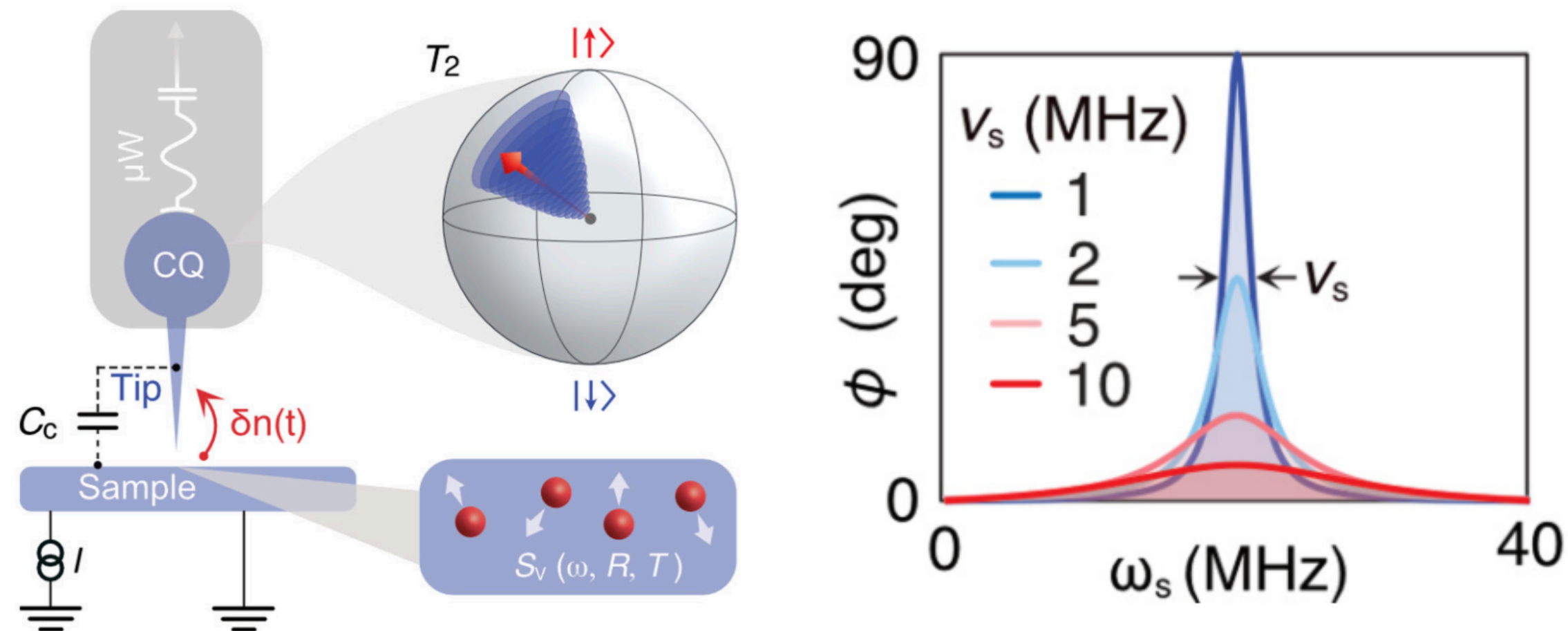
- Want to maximize  $C_s/C_{\text{tot}}$  so that the sample gates the qubit island effectively
- But capacitive coupling to a lossy sample limits sensitivity via  $T_1$
- Effect is small in CPB regime, but becomes significant as  $E_J/E_C$  is increased



# Measurement methods

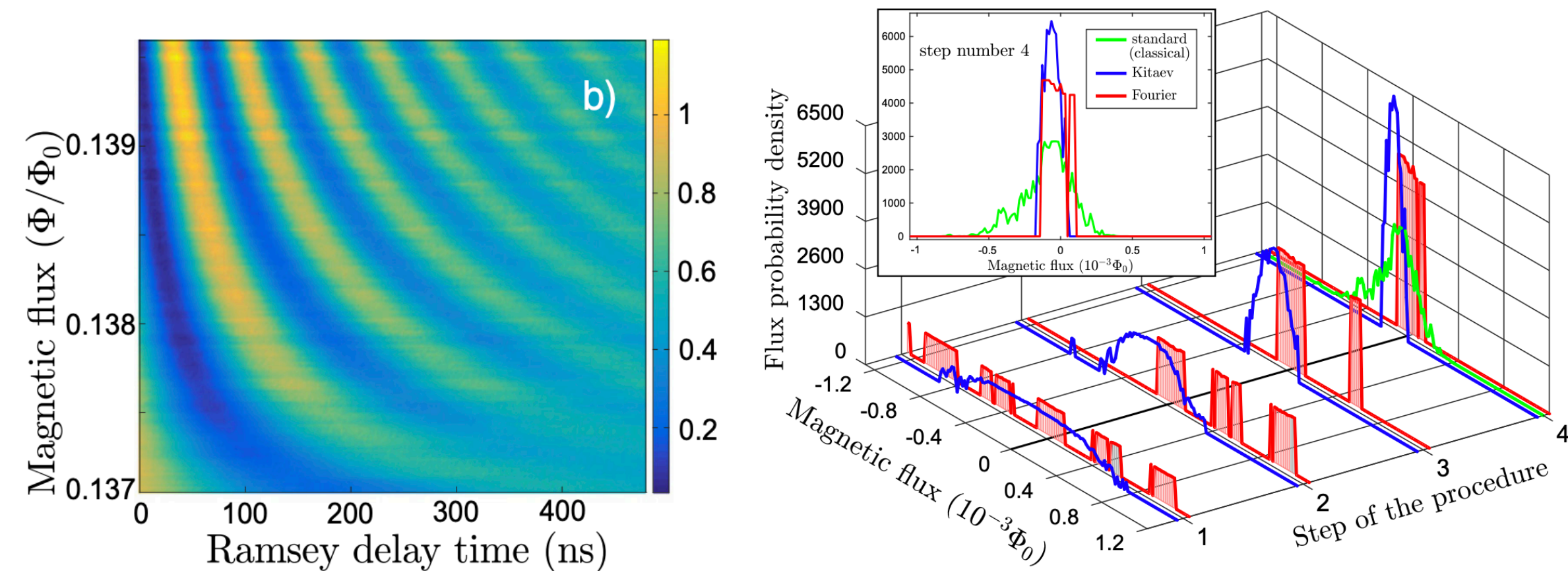
## Spectroscopy

- Can be done with only CW microwave control
- Center frequency measures the mean value of  $n_g$  or  $\Phi$
- Linewidth measures  $T_2$  but cannot easily be converted into a noise spectrum  $S_{n_g}(\omega)$  or  $S_{\Phi}(\omega)$
- Example: [Phys. Rev. Res. 2, 043031 \(2020\)](#)



## Interferometry

- Enables noise spectroscopy via dynamical decoupling
- Requires time domain control and reasonably high-fidelity readout
- Limited dynamic range due to phase wrapping
- Limited measurement repetition rate:
  - At most 1 bit of information per measurement, then the qubit must be reset to ground state
- Example: [npj Quant. Info. 4, 29 \(2018\)](#)



# Can we borrow ideas from SC qubit research without inheriting qubit downsides?

## Fast magnetic imaging/noise spectroscopy

- Requires mK temperatures
- Requires single-photon operation and few-photon readout
- Excess charge noise relative to SETs
- Limited measurement repetition rate
- Limited dynamic range

## SQUID magnetometer with dispersive readout

- Noise spectroscopy: high BW measurement + FFT
- Compatible with parametric amplifiers
- Faster measurement repetition rate (no reset)
- See:
  - [doi:10.1103/PhysRevB.83.134501](https://doi.org/10.1103/PhysRevB.83.134501)
  - [doi:10.1088/0953-2048/26/5/055015](https://doi.org/10.1088/0953-2048/26/5/055015)
  - [doi:10.1063/1.5030489](https://doi.org/10.1063/1.5030489)

## Fast potential imaging/noise spectroscopy

### Radio-frequency SET

- Noise spectroscopy: high BW measurement + FFT
- Lower charge noise than charge-sensitive qubit
- See:
  - [doi:10.1126/science.280.5367.1238](https://doi.org/10.1126/science.280.5367.1238)

## Imaging rf loss in SC circuits

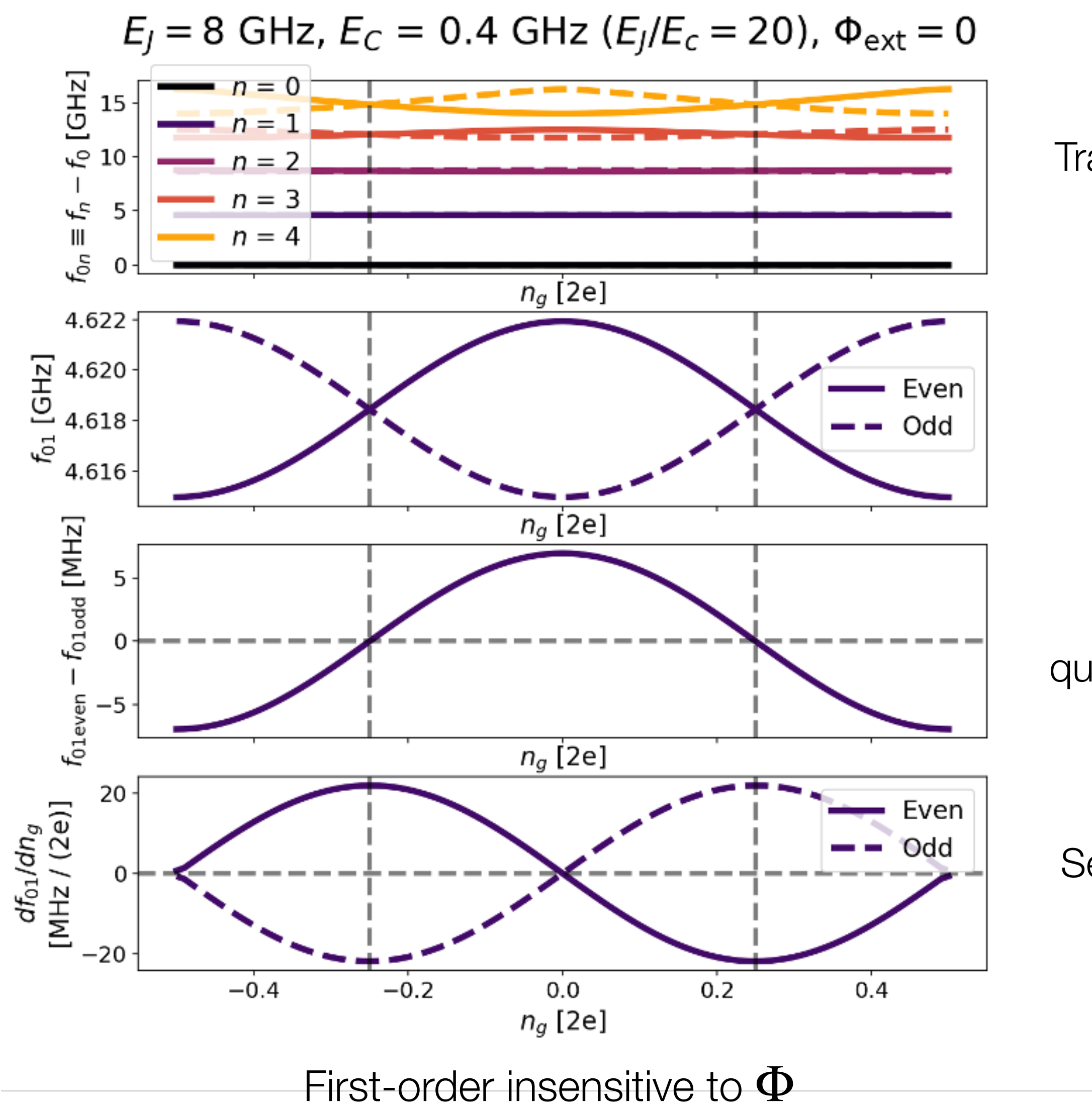
### Scanning high-Q SC resonator

- Measure Q vs. position
- Could be frequency-tunable with a SQUID
- See:
  - [doi:10.1063/1.4792381](https://doi.org/10.1063/1.4792381)
  - [doi:10.1103/PhysRevX.6.021044](https://doi.org/10.1103/PhysRevX.6.021044)





Goal: Measure qubit frequency  $f_{01}(n_g, \Phi)$  to quantify electric potential (via  $n_g$ ), magnetic field (via  $\Phi$ ), or fluctuations in those parameters (noise spectroscopy).

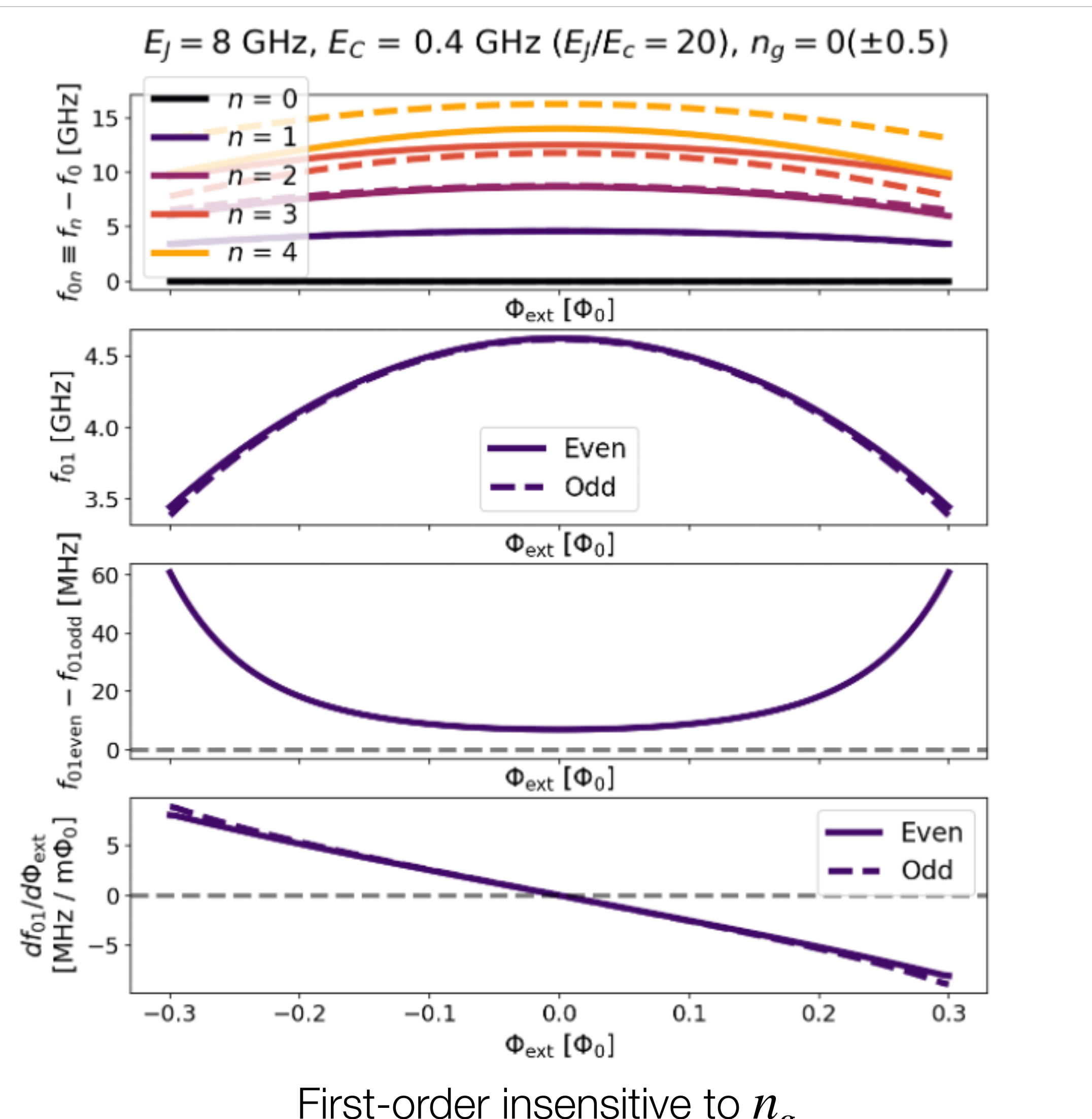


Transition frequencies

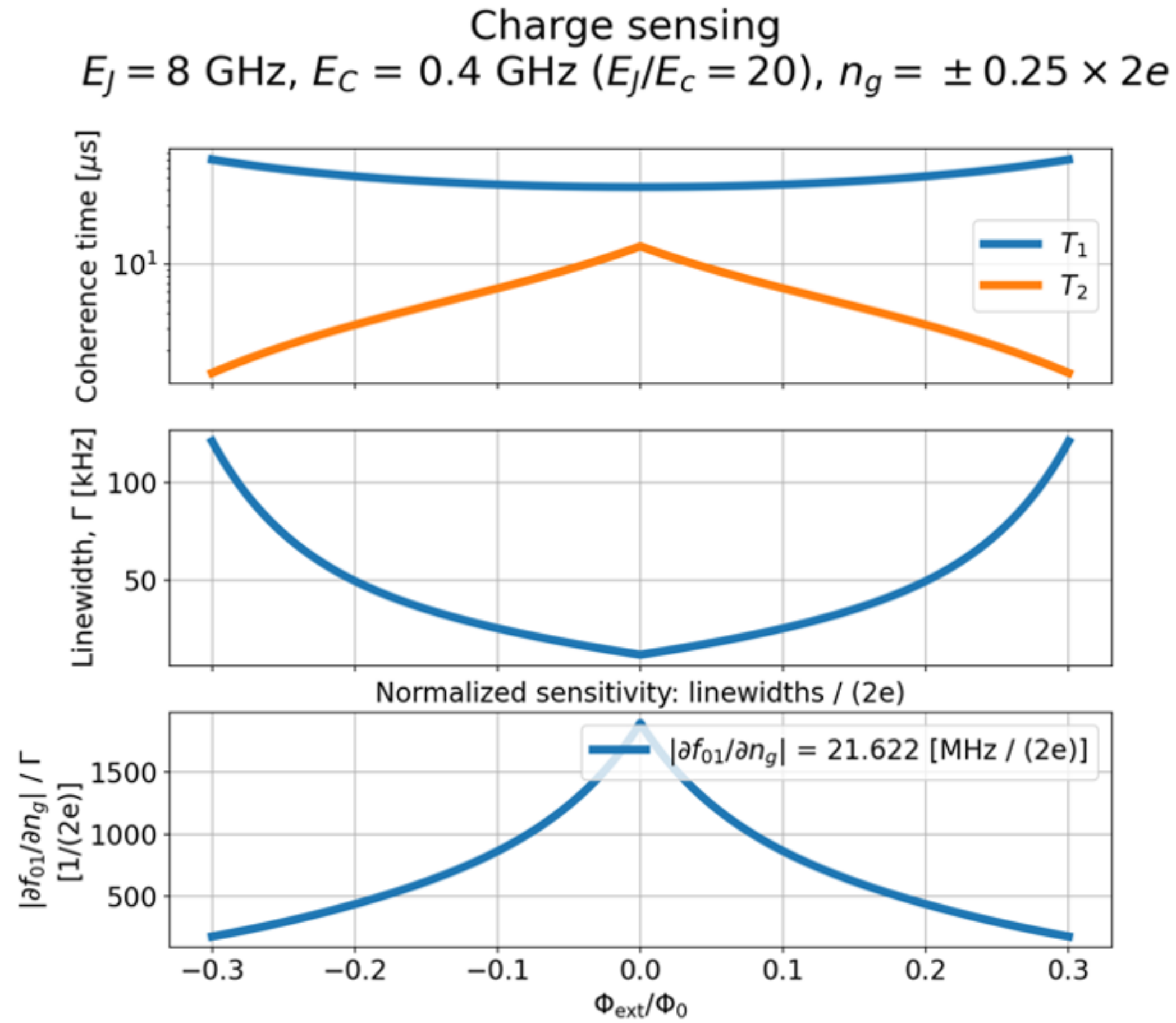
Qubit spectrum

Sensitivity to quasiparticle tunneling

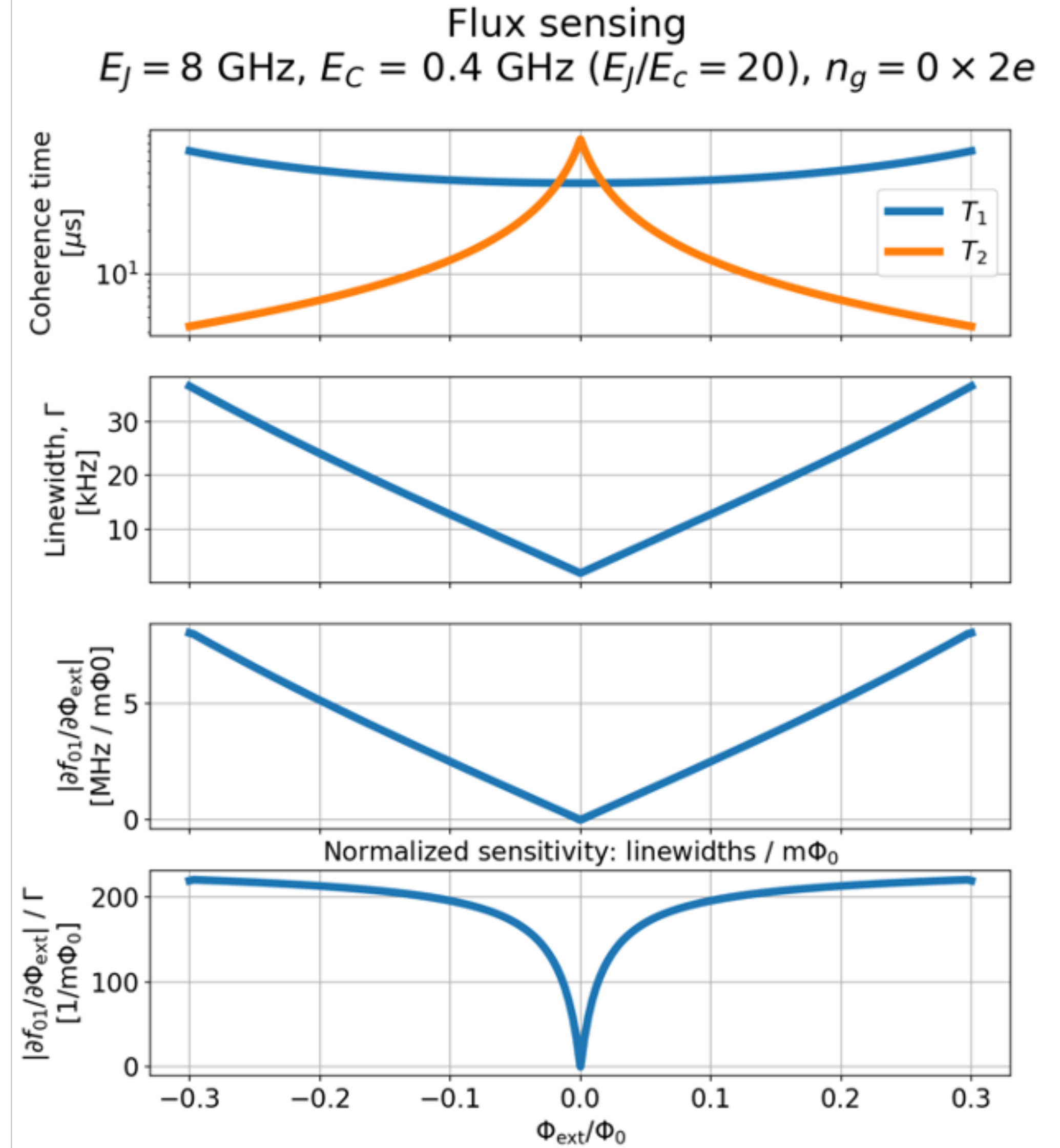
Sensitivity to  $n_g$  or  $\Phi$



$$A_{n_g} = 10^{-4} e / \sqrt{\text{Hz}}, \quad A_{\Phi} = 1 \mu\Phi_0 / \sqrt{\text{Hz}} \quad (@ 1 \text{ Hz})$$

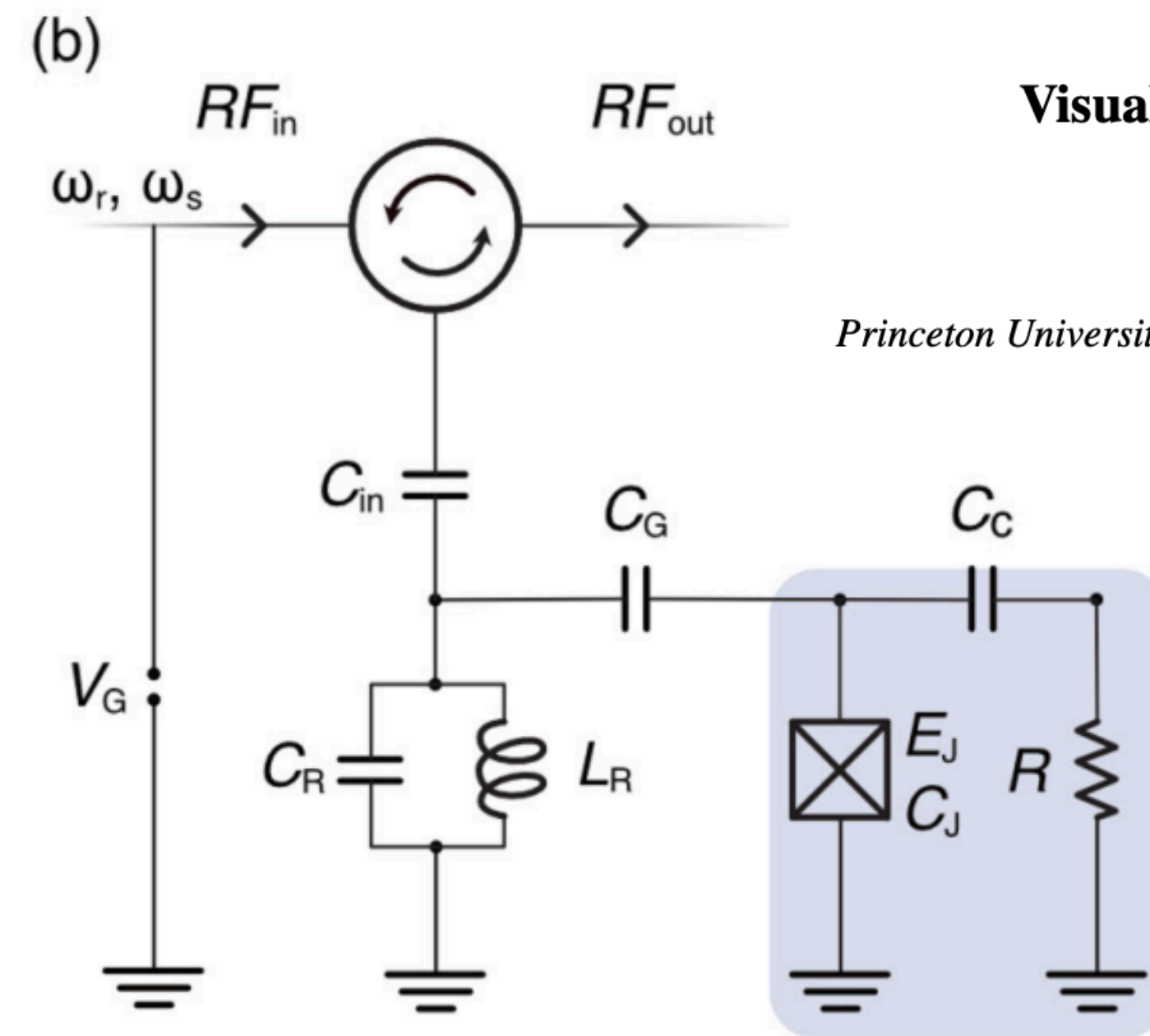
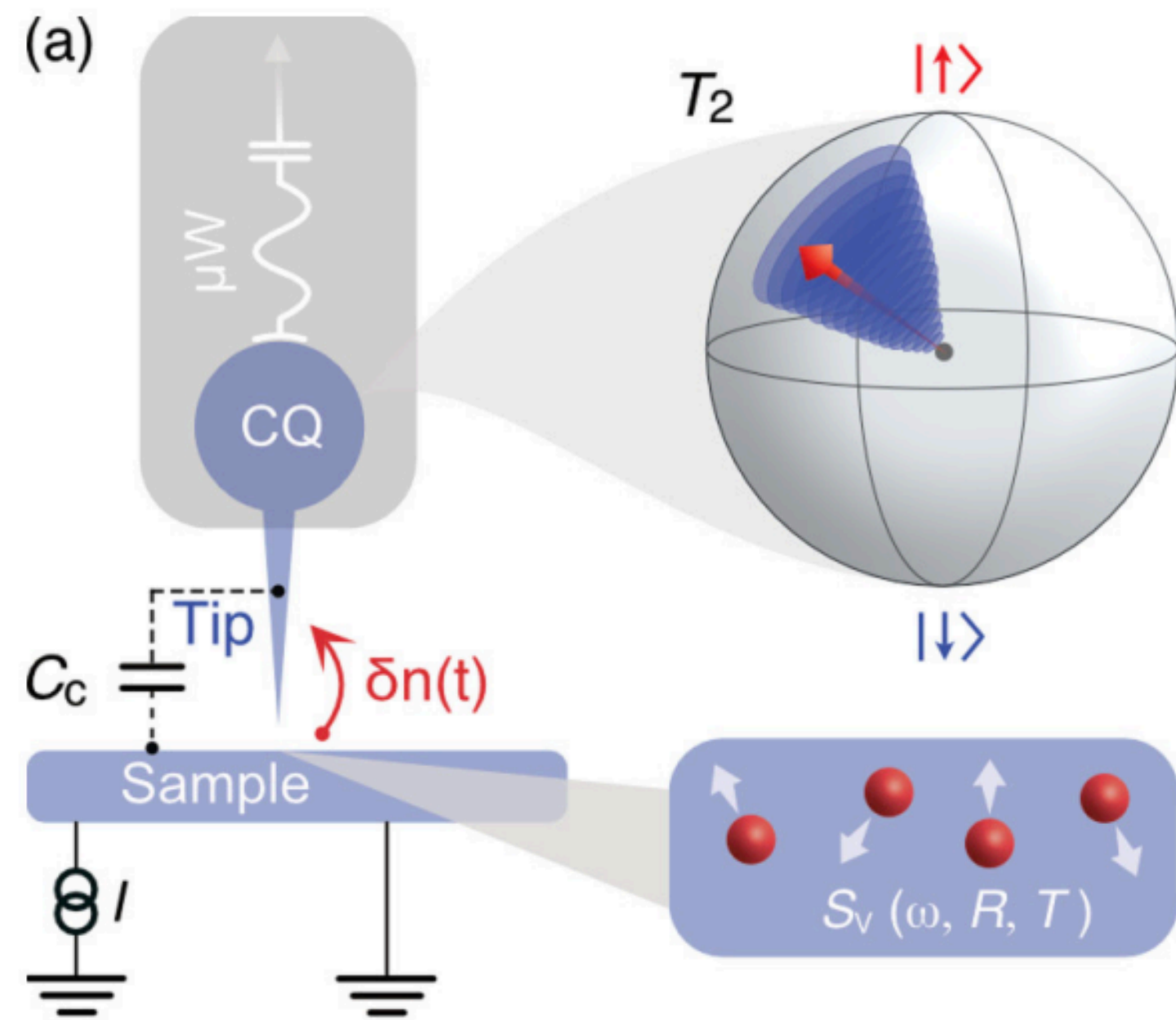


$$S_{n_g} = \frac{1}{\Gamma} \left| \frac{\partial f_{01}}{\partial n_g} \right| = 2\pi T_2 \left| \frac{\partial f_{01}}{\partial n_g} \right|$$



$$S_{\Phi} = \frac{1}{\Gamma} \left| \frac{\partial f_{01}}{\partial \Phi} \right| = 2\pi T_2 \left| \frac{\partial f_{01}}{\partial \Phi} \right|$$

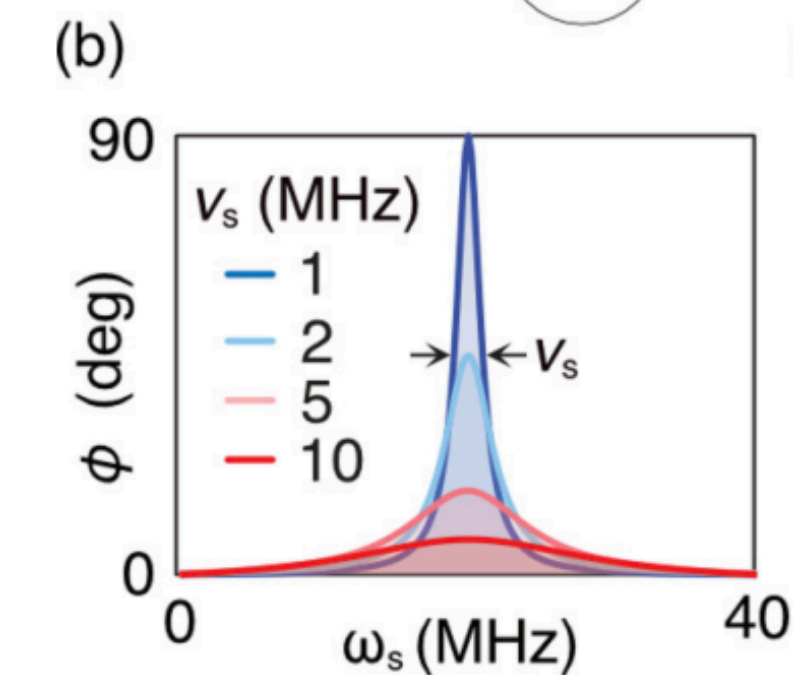
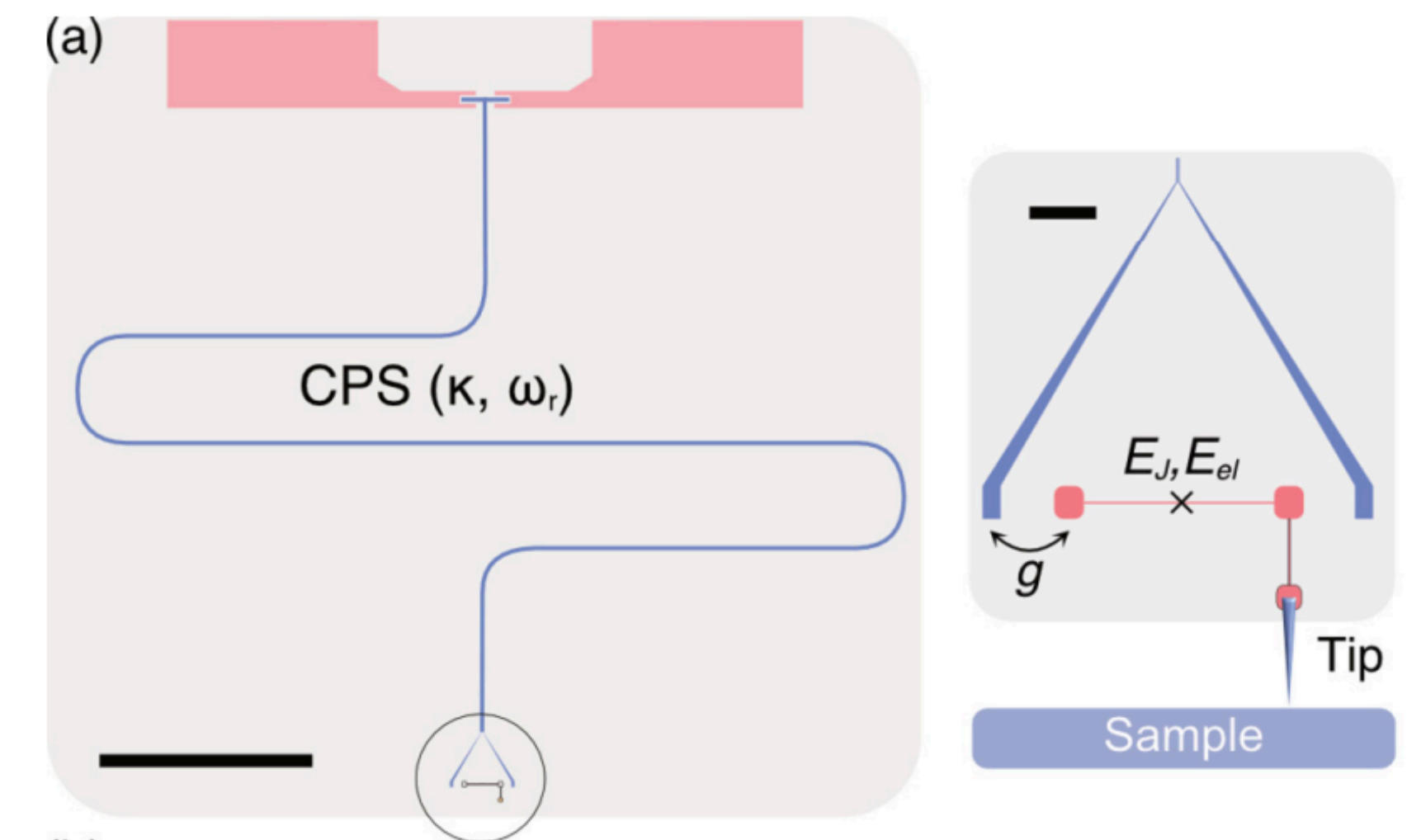




### Visualizing dissipative charge-carrier dynamics at the nanoscale with superconducting-charge-qubit microscopy

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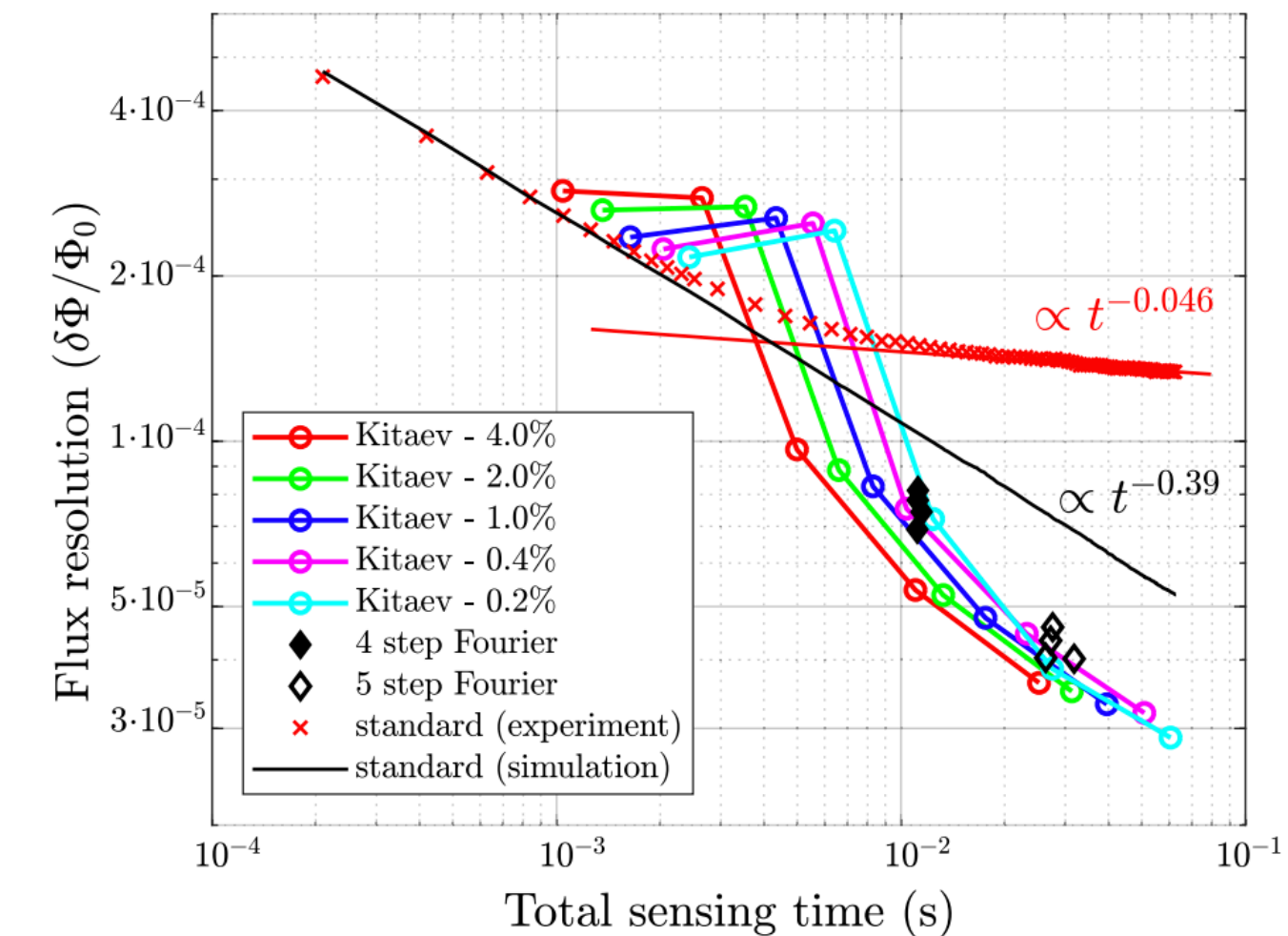
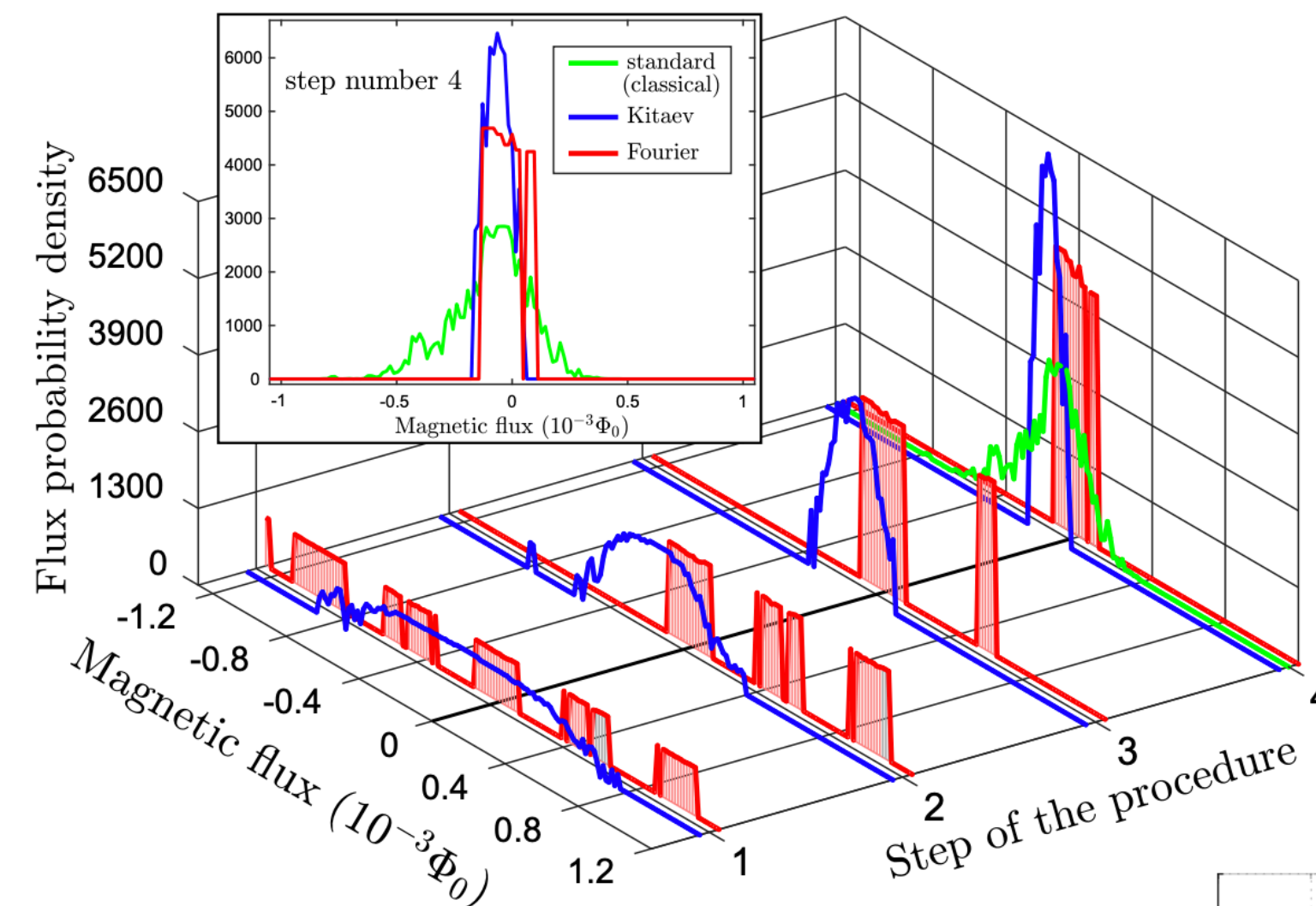
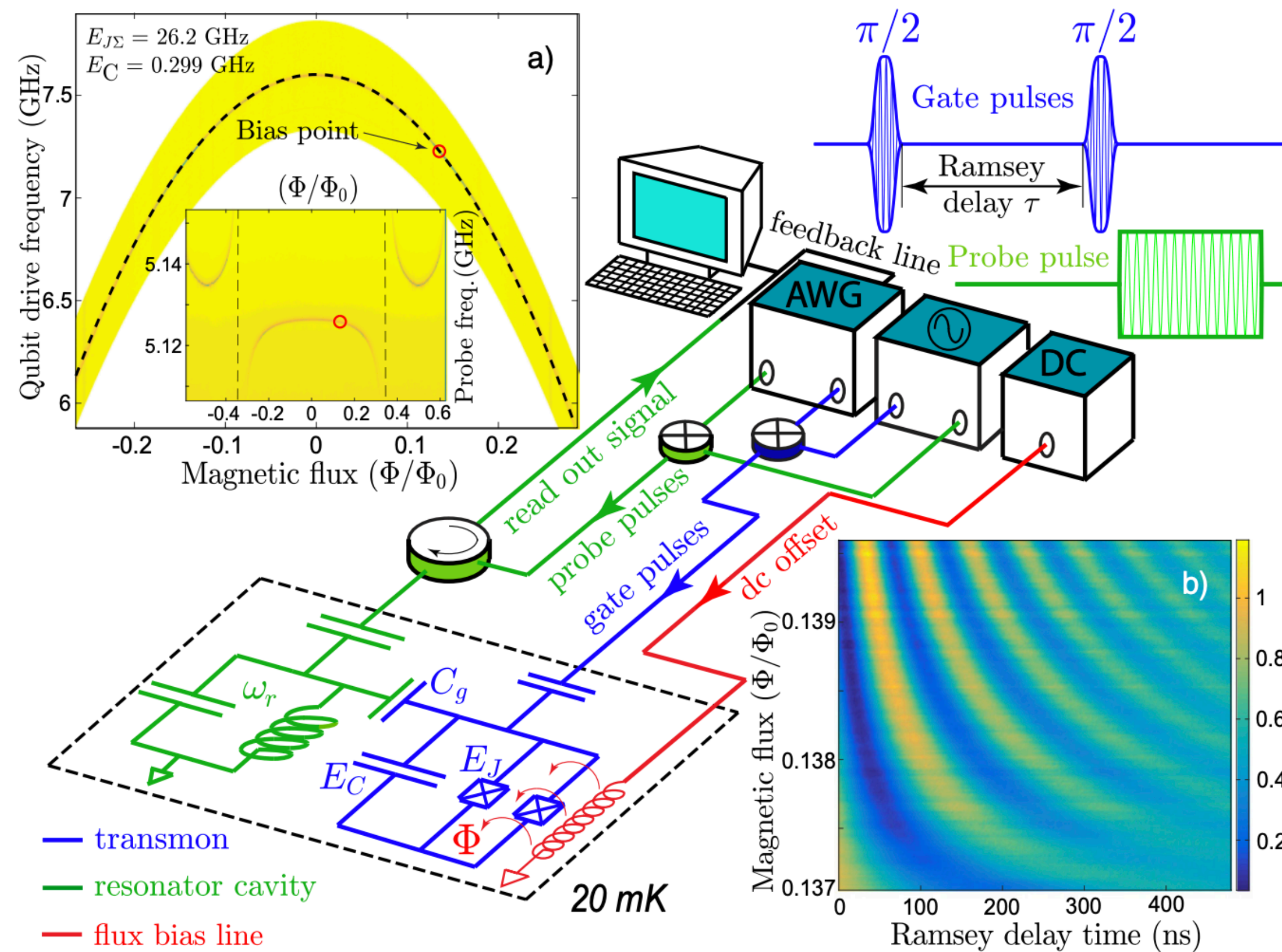
- Lossy sample capacitively coupled to a Cooper pair box
- Johnson noise from sample causes fluctuating  $n_g$ , reducing  $T_2$
- Measure linewidth  $1/T_2$  spectroscopically to infer local sample temperature, resistance



# Quantum-enhanced magnetometry by phase estimation algorithms with a single artificial atom

S. Danilin, A. V. Lebedev , A. Vepsäläinen, G. B. Lesovik, G. Blatter & G. S. Paraoanu 

*npj Quantum Information* **4**, Article number: 29 (2018) | [Cite this article](#)



- Tunable transmon,  $E_J/E_C \approx 90$
- Measure qubit frequency with a Ramsey sequence
- Use phase estimation algorithm to overcome dynamic range issue (phase wrapping)
- Sensitivity still limited by  $T_2$ , i.e.  $1/f$  flux noise



Qubit type	DC sensing	T2 noise spectroscopy (dynamical decoupling)	Scanning geometry	Advantages over [X]
<b>Single-junction CPB</b> $E_j/E_c \sim 1$	<b>Charge:</b> Spectroscopy only, limited by 1/f <b>Flux:</b> No	<b>Charge:</b> Low T2 makes time domain control difficult <b>Flux:</b> No	<b>Charge:</b> Seems doable <b>Flux:</b> --	<b>Charge (X = SET):</b> I don't see any <b>Flux:</b> --
<b>CPB with SQUID</b> $E_j/E_c \sim 1$	<b>Charge, flux:</b> Spectroscopy only, limited by 1/f, impractical measurement	<b>Charge, flux:</b> Low T2 makes time domain control difficult	<b>Charge + Flux:</b> Seems <u>pretty hard</u>	<b>Charge (X = SET):</b> I don't see any <b>Flux (X = SQUID):</b> I don't see any <b>Charge + Flux (X = SET + Hall):</b> I don't know
<b>Offset charge-sensitive transmon</b> $E_j/E_c \sim 10-20$	<b>Charge:</b> Spectroscopy or interferometry, limited by sample impedance or 1/f <b>Flux:</b> --	<b>Charge:</b> Yes, possibly limited by sample impedance <b>Flux:</b> --	<b>Charge:</b> Seems doable <b>Flux:</b> --	<b>Charge (X = SET):</b> Noise spectroscopy, can measure SC samples <b>Flux:</b> --
<b>Tunable transmon</b> $E_j/E_c \sim 50$	<b>Charge:</b> No <b>Flux:</b> Spectroscopy or interferometry	<b>Charge:</b> No <b>Flux:</b> Yes, requires fast control and readout	<b>Charge:</b> -- <b>Flux:</b> Seems doable	<b>Charge:</b> -- <b>Flux (X = SQUID):</b> Noise spectroscopy (is this more useful than a fast SQUID?)
<b>Flux qubit</b>	<b>Charge:</b> No <b>Flux:</b> Spectroscopy or interferometry	<b>Charge:</b> No <b>Flux:</b> Yes, requires fast control and readout	<b>Charge:</b> -- <b>Flux:</b> I don't know	<b>Charge:</b> -- <b>Flux (X = SQUID):</b> Noise spectroscopy (is this more useful than a fast SQUID?)